

# Design of a Miniaturized Superconducting Bandpass Filter by Evaluating the Kinetic Inductance in the $K$ -Inverter

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**SUMMARY** We propose a design theory of the miniaturized high temperature superconducting (HTS) coplanar waveguide (CPW) bandpass filter (BPF), which is composed of meanderline quarter-wavelength resonator,  $J$ - and  $K$ -inverters. The  $J$ - and  $K$ -inverters are realized by using interdigital gap and meander-shape inductor. To evaluate the kinetic inductance of the  $K$ -inverter, we fabricate the YBCO resonator connected with  $K$ -inverters and redesigned the YBCO filter parameters. Finally, we designed and fabricated the YBCO CPW quarter-wavelength resonator BPF by taking account of the kinetic inductance of the  $K$ -inverter. The experimental results are in agreement with the design parameters.

**key words:** HTS passive device, bandpass filter, quarter-wavelength resonator, meander line, coplanar waveguide

## 1. Introduction

High temperature superconducting (HTS) passive devices have an extremely high potential in the wireless base station systems, because of significant low insertion loss and high sensitivity. There are many reports on microstrip bandpass filter (BPF) [1]–[3], and HTS BPF based cryogenic receiver front-ends [4]. In order to reduce the cooling cost of the HTS filter subsystem, it is very important to miniaturize the HTS devices. A coplanar waveguide (CPW) is coated before patterning only one side of the substrate because the signal line and earth conductor exist only on the same side, and it is easy to make size reduction because the characteristic impedance of the transmission line depends only on the ratio of strip to slot width and does not depend on the thickness of the substrate. So, the CPW is more advantageous than the microstrip structure.

In our previous studies [5]–[7], we designed the miniaturized cross-coupled CPW BPF by using highly packed meanderline half-wavelength ( $\lambda/2$ ) resonators and interdigital gaps ( $J$ -inverters). In this paper, for further miniaturization, we designed and tested quarter wavelength ( $\lambda/4$ ) resonator BPF by using interdigital gaps and meander-shape inductors ( $K$ -inverters). However, the effect of the kinetic inductance of HTS appears conspicuously and quality factor becomes lower because the microwave current concentrates on the narrow meander inductor. The effect of the kinetic inductance changes the bandwidth of the BPF (see Eqs. (3)), and lower quality factor causes the insertion loss

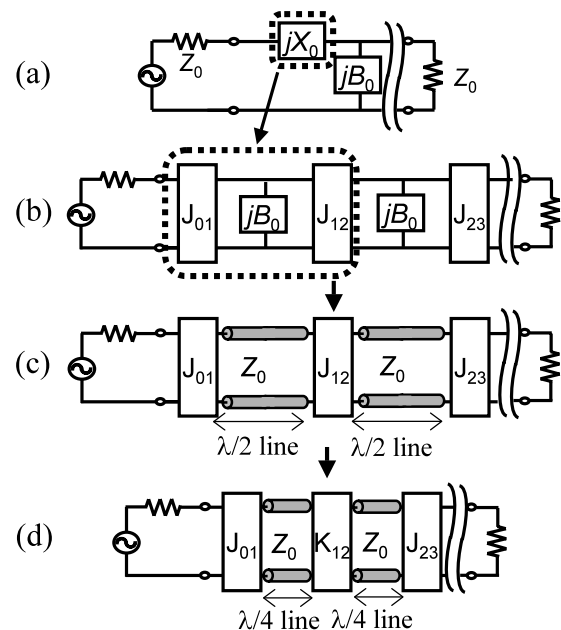
in the pass band. However, most important problem for the narrow band HTS BPF is the difference of the bandwidth between the design parameter and experimental result.

In order to take account of the kinetic inductance in the  $K$ -inverter, we fabricate the YBCO resonator connected with  $K$ -inverters and redesigned the YBCO filter parameters by using the experimental results. The prototype YBCO miniaturized BPF (center frequency = 10 GHz, 2-pole, bandwidth=1.5%), which includes the kinetic inductance of the  $K$ -inverter, is also tested in cryogenic temperature. We show that the experimental results are in agreement with the design parameters.

## 2. Design of Chebyshev Quarter-Wavelength Resonator CPW Bandpass Filter

### 2.1 Design Parameters of the $\lambda/4$ Resonator BPF

Figure 1(a) shows the lumped element circuit model of the BPF [8]. Series resonator ( $jX_0$ ) is transformed into the shunt resonator ( $jB_0$ ) by interconnecting two  $J$ -inverters as shown in Fig. 1(b). Figure 1(c) and Fig. 1(d) shows the equivalent



**Fig. 1** Lumped element circuit model of the BPF (a), circuit model composed of the shunt resonators and  $J$ -inverters (b), equivalent circuit model of the BPF with  $\lambda/2$  resonators (b) and  $\lambda/4$  resonators (c).

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circuit model of the BPF with  $\lambda/2$  resonators and  $\lambda/4$  resonators, respectively [9]. In Fig. 1(d), the  $\lambda/4$  resonator is connected with the  $J$ - and  $K$ -inverters, so that total length of the BPF is almost as half as that of the  $\lambda/2$  resonator BPF.

The design parameters of the Chebyshev type  $\lambda/2$  resonator BPF are given by [8]

$$\begin{cases} J_{01} = \sqrt{w} \sqrt{\frac{Y_0 b_0}{g_0 g_1}} \\ J_{i,i+1} = w \sqrt{\frac{b_0 b_0}{g_i g_{i+1}}} \quad (i = 1, 2, \dots, n-1) \\ J_{n,n+1} = \sqrt{w} \sqrt{\frac{Y_0 b_0}{g_n g_{n+1}}} \\ b_0 = \frac{\omega_0}{2} \left. \frac{\partial B_0}{\partial \omega} \right|_{\omega=\omega_0} = \frac{\pi}{4} Y_0 = \frac{\pi}{4} \frac{1}{Z_0} \end{cases}, \quad (1)$$

Where,  $w$  and  $g_i$  are the fractional bandwidth and element values of Chebyshev BPF, respectively.  $b_0$  is the susceptance slope parameter of the shunt resonator ( $jB_0$ ) at the resonant frequency ( $\omega_0$ ) (see Fig. 1(a)).

In order to realize the miniaturized BPF, these resonators are modified and bended to the meanderline structures. The bended resonators change the slope parameters because of the interconnection between the resonators. The reactance slope parameter ( $x_i$ ) and susceptance slope parameter ( $b_j$ ) of the  $i$ -th resonator ( $jX_i$ ) and  $j$ -th resonator ( $jB_j$ ) can be calculated and given by

$$\begin{cases} x_i = \frac{\omega_0}{2} \left. \frac{\partial X_i}{\partial \omega} \right|_{\omega=\omega_0} = \frac{\pi}{4} Z_i \\ b_j = \frac{\omega_0}{2} \left. \frac{\partial B_j}{\partial \omega} \right|_{\omega=\omega_0} = \frac{\pi}{4} Y_j \end{cases}, \quad (2)$$

where,  $Z_i$  and  $Y_j = (1/Z_j)$  are characteristic impedance and characteristic admittance, respectively.

The design parameters of the even pole number  $\lambda/4$  resonator BPF are given by

$$\begin{cases} J_{01} = \sqrt{w} \sqrt{\frac{Y_0 b_1}{g_0 g_1}} = \frac{\pi}{4} \sqrt{w} \sqrt{\frac{1}{g_0 g_1 x_1 Z_0}} \\ K_{i,i+1} = w \sqrt{\frac{x_i x_{i+1}}{g_i g_{i+1}}} = \left(\frac{\pi}{4}\right)^2 w \sqrt{\frac{1}{g_i g_{i+1} b_i b_{i+1}}} \\ \quad (i = 1, 3, \dots, 2n-1) \\ J_{i,i+1} = w \sqrt{\frac{b_i b_{i+1}}{g_i g_{i+1}}} = \left(\frac{\pi}{4}\right)^2 w \sqrt{\frac{1}{g_i g_{i+1} x_i x_{i+1}}} \\ \quad (i = 2, 4, \dots, 2n-2) \\ J_{n,n+1} = \sqrt{w} \sqrt{\frac{Y_0 b_n}{g_n g_{n+1}}} = \frac{\pi}{4} \sqrt{w} \sqrt{\frac{1}{g_n g_{n+1} x_n Z_0}} \end{cases}. \quad (3)$$

We can see  $K$ - and  $J$ -inverters change the bandwidth ( $w$ ) of the BPF.

## 2.2 Design of the $K$ -Inverter and $J$ -Inverter

The exact inverters, and slope parameters of the  $\lambda/4$  resonators are calculated by using the electromagnetic field

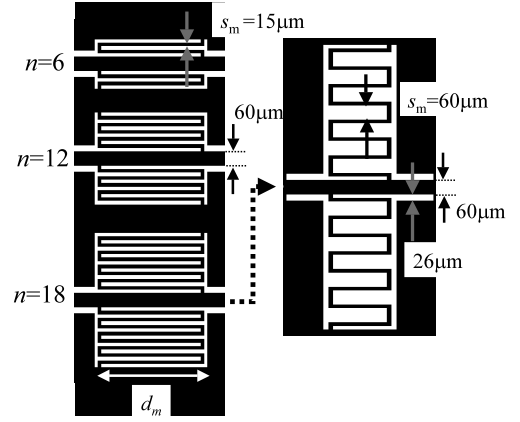


Fig. 2 Layout of the meander inductor ( $K$ -inverter).  $d_m$ ,  $s_m$  and  $n$  is the length, gap width and bending numbers of the  $K$ -inverters, respectively.

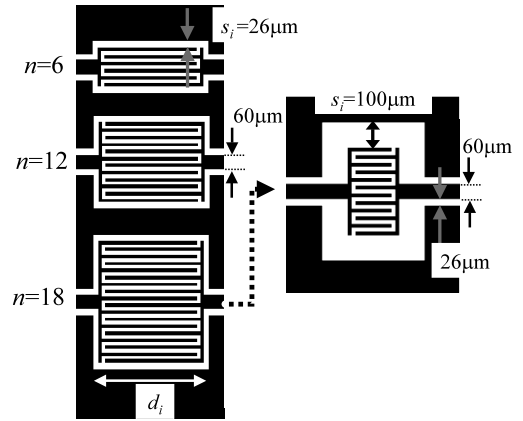
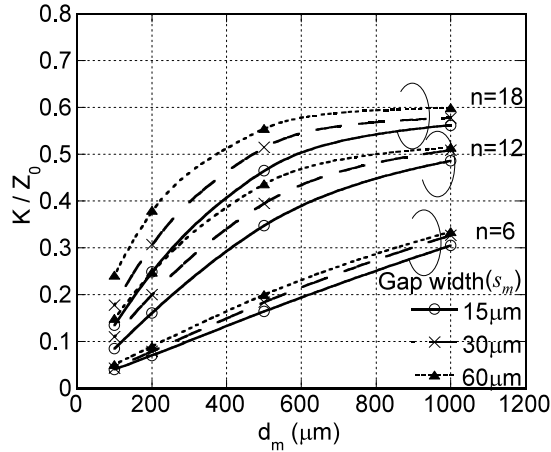


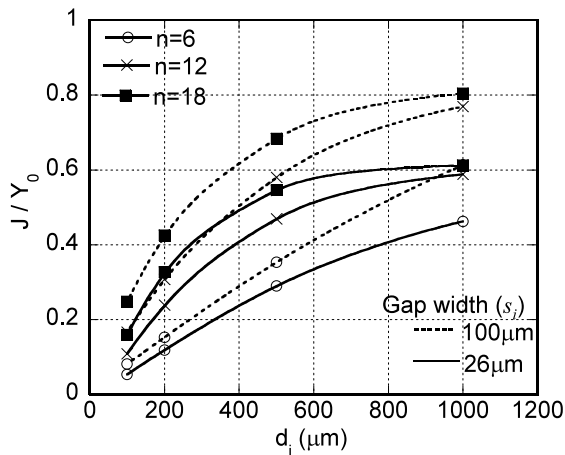
Fig. 3 Layout of the interdigital capacitor ( $J$ -inverter).  $d_i$ ,  $s_i$  and  $n$  is the length, gap width and bending numbers of the  $J$ -inverters, respectively.

simulator (EM-simulator) (Agilent, ADS2004A). Figure 2 shows the layout of the  $K$ -inverters. They are realized by using meanderline inductors. The YBCO film is approximated as the perfect conductor and the thickness of the MgO ( $\epsilon_r=9.6$ ) substrate is  $500\mu\text{m}$ . The width of the signal line and gap interval is  $60\mu\text{m}$  and  $26\mu\text{m}$  (characteristic impedance  $=50\Omega$ ), respectively (see Fig. 2). In the figure,  $d_m$  and  $n$  is the length and bending number of the  $K$ -inverters. In the previous paper, we found that  $K$  values are saturated even  $d_m$  becomes longer [9]. In this paper, in order to design the BPFs, which have a lot of pole number, fractional bandwidth and so on, we modified the shape of the inverters, which is adjustable in wide range. In order to suppress the interference of the magnetic field around the meander inductor, the gap ( $s_m$ ) is stretched from  $15\mu\text{m}$ , to  $60\mu\text{m}$  (see Fig. 2).

Figure 3 shows the layout of the  $J$ -inverters. In the figure,  $d_i$  and  $n$  is the length and finger numbers of the  $J$ -inverters, respectively. We also confirmed that  $J$  values are saturated even  $d_i$  becomes longer [9]. In order to suppress the effect of the shunt capacitance between signal line and earth conductor, the gap ( $s_i$ ) is stretched from  $26\mu\text{m}$ , to



**Fig. 4** EM-simulation results of the  $d_m$  and  $s_m$  dependences on  $K$ -inverters.



**Fig. 5** EM-simulation results of the  $d_i$  and  $s_i$  dependences on  $J$ -inverters.

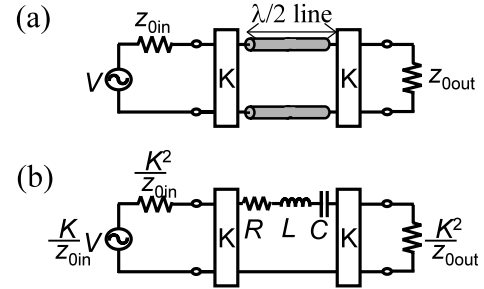
100  $\mu\text{m}$  (see Fig. 3).

Figure 4 shows the EM-simulation results of the  $d_m$  and  $s_m$  dependence on  $K$ -inverters.  $K$  values are normalized by  $Z_0$  ( $=50 \Omega$ ).  $K/Z_0$  are both increased by stretching the gap, but are saturating in large  $d_m$  regions. Although large  $K$ -values are obtained in large  $s_m$  region, the chip area increases. Also, the  $J/Y_0$  are simulated and measured in the same method, where  $Y_0$  ( $=1/Z_0$ ) is 0.02S.

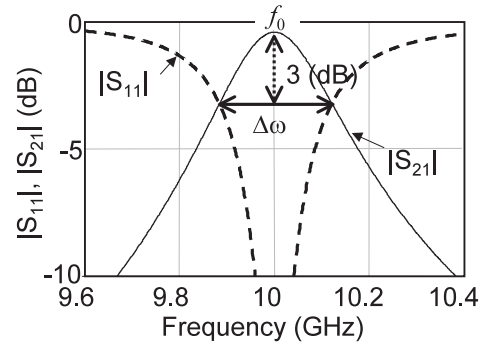
Figure 5 shows the EM-simulation results of the  $d_i$  and  $s_i$  dependence on  $J$ -inverters.

### 2.3 Design of the $K$ -Inverter by Evaluating the Kinetic Inductance in the YBCO Film

In order to investigate the kinetic inductance of the  $K$ -inverter, we fabricate the YBCO  $\lambda/2$  resonator connected with  $K$ -inverters and measure these effects. Figure 6(a) shows the transmission line model of the  $\lambda/2$  resonator with  $K$ -inverters.  $Z_{0in}$  and  $Z_{0out}$  are source and load impedance, respectively. Because  $K$ -inverters are connected to the ground plane,  $\lambda/2$  transmission line is transformed into the



**Fig. 6**  $\lambda/2$  transmission line model (a) and equivalent circuit model of the resonator (b) connected with the  $K$ -inverters.



**Fig. 7** Typical transmission characteristics of the resonator for  $K$  value estimation (EM-simulation results).

series resonator ( $R$ ,  $L$  and  $C$ ) at resonant frequency as shown in Fig. 6(b).

Figure 7 shows the EM-simulation results of typical transmission characteristics of the resonator for  $K$  value estimation. As shown in Fig. 6(b), input and output impedances are inverted with  $K$ -inverters.

An available power from the source ( $P_{av}$ ) is given by,

$$P_{av} = \frac{|V|^2}{8Z_{0in}}. \quad (4)$$

When a power delivered to the load ( $Z_{0out}$ ) is  $P_{out}$ , the ratio of the power ( $P_{out}/P_{av}$ ) at  $f$  ( $=\omega/2\pi$ ) is given by

$$\frac{P_{out}}{P_{av}} = |S_{21}(\omega)|^2 = \frac{\left. \frac{P_{out}}{P_{av}} \right|_{\omega=\omega_0}}{1 + Q_L^2 \left( \frac{\omega}{\omega_0} - \frac{\omega_0}{\omega} \right)^2}, \quad (5)$$

$$\begin{aligned} \left. \frac{P_{out}}{P_{av}} \right|_{\omega=\omega_0} &= |S_{21}(\omega_0)|^2 \\ &= \frac{4 \frac{K^2}{Z_{0in}} \frac{K^2}{Z_{0out}}}{\left( \frac{K^2}{Z_{0in}} + R + \frac{K^2}{Z_{0out}} \right)^2} = \frac{4Q_L^2}{Q_{e1}Q_{e2}}, \end{aligned} \quad (6)$$

where, the loaded quality factor ( $Q_L$ ) is calculated from the unloaded quality factor ( $Q_u$ ) and external quality factor ( $Q_{e1}$  and  $Q_{e2}$ ), and are given by

$$Q_L \equiv \frac{1}{1/Q_{e1} + 1/Q_u + 1/Q_{e2}} = \frac{x}{K^2/Z_{0in} + R + K^2/Z_{0out}}, \quad (7)$$

$$Q_u \equiv \frac{\omega_0 L}{R} = \frac{x}{R}, \quad (8)$$

$$Q_{e1} \equiv \frac{\omega_0 L}{K^2/Z_{0in}} = \frac{x}{K^2/Z_{0in}}, \quad (9)$$

$$Q_{e2} \equiv \frac{\omega_0 L}{K^2/Z_{0out}} = \frac{x}{K^2/Z_{0out}}. \quad (10)$$

When  $Z_{0in} = Z_{0out} = Z_0$ , we can approximate  $Q_{e1} = Q_{e2} = Q_e$ . Finally,  $Q_L$ ,  $Q_u$ ,  $Q_e$  and  $K$  are given by

$$Q_L = \frac{\omega_0}{\Delta\omega}, \quad (11)$$

$$Q_u = \frac{Q_L}{1 - |S_{21}(\omega_0)|}, \quad (12)$$

$$Q_e = \frac{2Q_L}{|S_{21}(\omega_0)|}, \quad (13)$$

$$K = \sqrt{\frac{Z_0\pi}{2Q_e}}. \quad (14)$$

We can calculate the  $K$  value from the transmission coefficient ( $|S_{21}|$ ).

### 3. Results and Discussion

#### 3.1 Experimental Results of $K$ -Inverter

Figure 8 shows the photograph of the YBCO CPW  $\lambda/2$  resonator connected with  $K$ -inverters. In order to improve the accuracy of the RF signals, the air coplanar probes (Cascade Microtech, GSG-150) touched the YBCO transmission line (see Fig. 8). The width of the signal line and gap interval is  $60\mu\text{m}$  and  $26\mu\text{m}$ . The meander interval of the  $\lambda/2$  resonator is  $300\mu\text{m}$ . Figure 9 shows the photograph of the  $K$ -inverter. There were fabricated 2 types of the meander

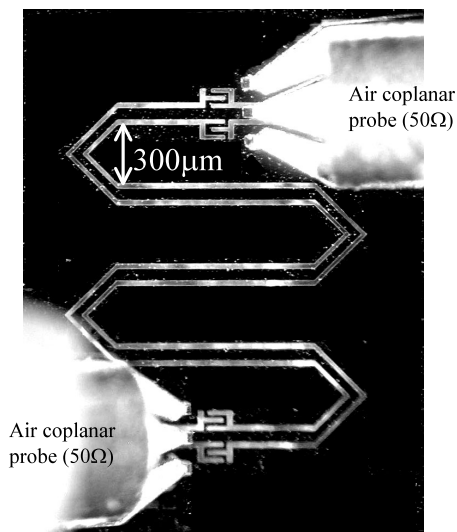


Fig. 8 Photograph of the YBCO CPW  $\lambda/2$  resonator with probs.

inductor ( $n=6$ ,  $d_m=150\mu\text{m}$  and  $200\mu\text{m}$ ). These resonators are placed in a vacuum chamber with refrigerator cooling system. We measured the  $S$  parameters by a GP-IB controlled vector network analyzer (HP, HP8722C).

Figure 10 shows the frequency response of the  $|S_{21}|$  of the 2 resonators named #1 and #2 at 20 K. We can calculate the  $K$  values by using Eqs. (11)–(14). Table 1 shows the normalized  $K$  values from the experimental results, where  $Z_0=50\Omega$ . Experimental results are approximately 50% as large as that of the simulation results. This EM simulator cannot evaluate the kinetic inductance ( $L_k$ ) of the YBCO meander inductor, namely it is considered only a magnetic inductance ( $L_m$ ). However, the experimental results ( $L_{\text{exp}}$ ) include  $L_k$  and given by,

$$L_{\text{exp}} = L_m + L_k. \quad (15)$$

It is considered that the effect of the kinetic inductance of HTS appears conspicuously because the microwave current concentrates on the thin, narrow and long meander inductor ( $K$ -inverter section). Because  $L_m/L_k$  increases during the

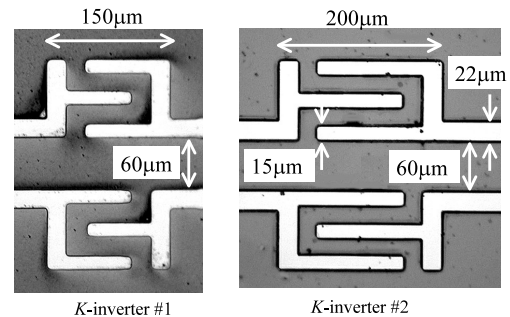


Fig. 9 Photograph of the  $K$ -inverters.

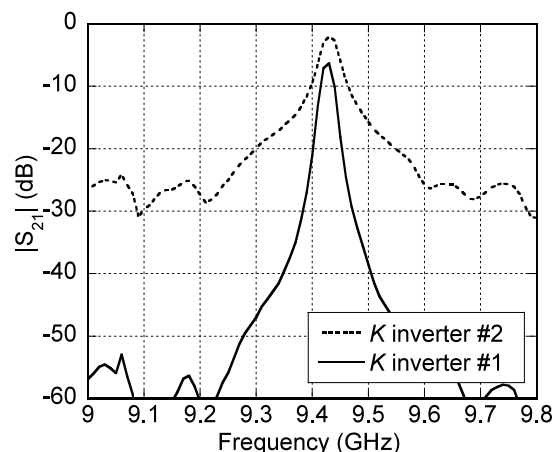


Fig. 10 Transmission characteristics of the resonators at 20 K.

Table 1 Normalized  $K$  values of the inverters.

K inverters	$K/Z_0$ (Sim.) (A)	$K/Z_0$ (Exp.) (B)	A/B (%)
$K$ inverter #1	0.0291	0.0667	43.64
$K$ inverter #2	0.0358	0.0701	50.98

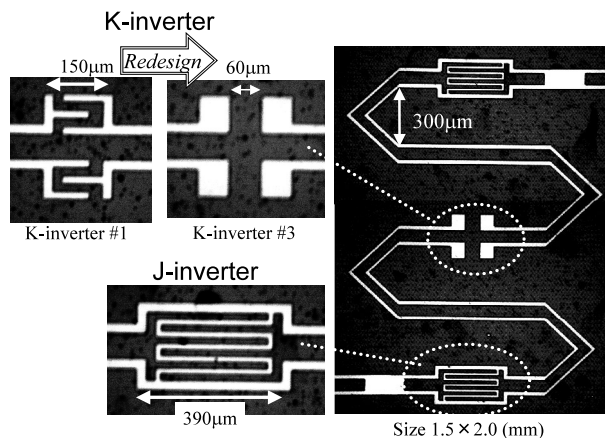


Fig. 11 Photograph of the prototype YBCO BPF with CPW  $\lambda/4$  resonator (center frequency = 10 GHz, 2-pole, bandwidth = 1.5%).

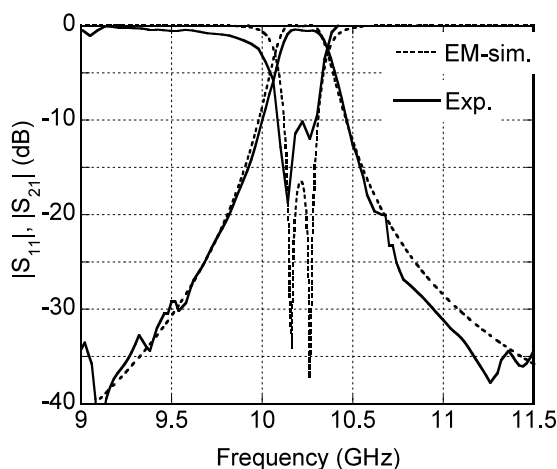


Fig. 12 Frequency response of the prototype YBCO CPW  $\lambda/4$  resonator BPF at 20 K.

miniaturization, and  $L_k$  changes by the geometry factor, it is necessary to obtain the  $L_{exp}$  for fabricating  $\lambda/4$  resonator BPF.

### 3.2 Experimental Results of a Prototype YBCO BPF

Figure 11 shows the photograph of the prototype YBCO BPF with  $\lambda/4$  resonator for 10 GHz application (2-pole, bandwidth=1.5%). Using the experimental results of the inverters, we redesign the shape of the inverters by evaluating the kinetic inductance. In the figure, the meander inductor and interdigital gap are also shown. The filter size is  $1.5 \times 2.0$  mm. We also modified the shape of the  $K$ -inverter as above process. The desirable  $K$  value is  $K/Z_0 = 0.0291$  ( $d_m = 150 \mu\text{m}$ ), which is realized by the  $K$ -inverter #1 in the EM simulator (without evaluating the  $L_k$ ), however, actual  $K$  value is realized by  $K$ -inverter #3 which has not the meander shape as shown in Fig. 11.

Figure 12 shows the frequency responses of the 2-pole YBCO CPW  $\lambda/4$  BPF at 20 K. In the figure, broken lines show the  $S$  parameters of the EM simulation results. Be-

cause we evaluate the actual  $K$  values of the inverter experimentally, and redesign and fabricate the BPF, the bandwidth is almost similar to that of the EM simulation results.

## 4. Conclusion

$K$ -inverter, which composed of a meander-shape inductor, is effective for the HTS microwave devices because of the extremely high quality factor. However, the effect of the kinetic inductance increases, as HTS BPF has thin, narrow and long transmission lines. We can design and fabricate the miniaturized YBCO CPW quarter-wavelength resonator BPF by experimentally evaluating the kinetic inductance of the  $K$ -inverter.

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